

# Field-Extracted Lumped-Element Models of Coplanar Stripline Circuits and Discontinuities for Accurate Radio-Frequency Design and Optimization

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**Abstract**—In this paper, we present original lumped-element models of a large variety of coplanar stripline (CPS) circuits and discontinuities for accurate design and optimization of radio-frequency integrated circuits. These circuit models are extracted by applying a recently proposed deembedding technique called short-open calibration to calibrate the calculated field parameters obtained from the full-wave method of moments (MoM). It is realized by defining the two calibration standards, i.e., short and open elements, to evaluate and remove the error terms existing in the deterministic MoM algorithm, such as the approximation of port discontinuity and inconsistency of two-dimensional and three-dimensional characterizations of CPS external lines. In contrast to the static models, these field-extracted models can account for all the physical effects subject to the core area of CPS discontinuity, including frequency dispersion, high-order modes, and radiation loss. With this scheme, several CPS circuits and discontinuities are investigated over a wide frequency range to formulate their lumped-element models and also to expose their interesting features in connection with electrically finite width of the two strip conductors.

**Index Terms**—Coplanar stripline discontinuity, lumped-element model, method of moments, short-open calibration technique.

## I. INTRODUCTION

COPLANAR stripline (CPS) [1] has recently gained a significant momentum in the design and applications of high-density radio-frequency integrated circuits (RFICs). The CPS geometry may be viewed as the dual structure of coplanar waveguide (CPW). With the arrangement of its two strips on a common dielectric interface, the CPS has all of the attractive features of its CPW counterpart, namely, easy insertion of shunt and series passive and active devices, and no need for via-holes to connect the ground plane [2]. In addition, the CPS presents some additional advantages over the CPW in that it eliminates the requirement of designing air-bridges around discontinuity to suppress the unwanted mode, as in the CPW

[2]. The balanced geometry of the CPS offers flexibility in the design of uniplanar balanced circuits such as a uniplanar broad-band double-balanced mixer [3] and lumped-element CPS low-pass filter [4], as well as a CPS-fed printed dipole antenna array [5].

As compared to the extensive work on microstrip and CPW structures, few investigations have been carried out to date for studying electrical behavior of CPS discontinuities/elements and developing their computer-aided design (CAD) models that are useful for the design of passive and active CPS circuits and antennas. Very recently, a limited characterization of some simple CPS circuits [6]–[8] has been reported with resorting to a finite-difference time-domain (FDTD) method, as well as experimental procedures. Otherwise, a full-wave spectral-domain method was used for the characterization of a simple CPS open circuit [9], [10]. Still, it is desirable to formulate accurate equivalent lumped-element models of a large variety of CPS basic circuits and discontinuities for general-purpose design and optimization. Parasitic and dispersion effects should be rigorously accounted for in the models, which can be achieved only by a parameter extraction from some adequate field-based technique. This is important for a uniplanar CPS structure because its ground plane is always finite and it may generate intolerable parasitic effects, as will be shown in this paper. Specifically, additional series inductance and potential radiation caused by some CPS discontinuities in unbalanced CPS circuits, namely, CPS gap, may be quite severe, which cannot be ignored in the models.

In response to the above demands, a deterministic method of moments (MoM) algorithm, as detailed in [11], is extended for field modeling of the CPS structures. This is made possible by introducing an impressed electric field across the two strip conductors of CPS feed lines at the port planes of choice [12], [13]. Inspired from the concept of experimental calibration procedure [14], a numerical calibration technique, namely, short-open calibration (SOC) [15]–[17], was implemented to calibrate the MoM-based calculated parameters by removing the error terms existing in the 3-D MoM algorithm. These terms stand for the approximation of impressed fields at the ports and inconsistency of two-dimensional (2-D) and three-dimensional (3-D) characterizations of CPS external lines, when the deembedding procedure is required in the formulation of equivalent CAD model as to the preferable reference planes. With the proposed SOC

Manuscript received May 26, 1999; revised November 6, 2000. This work was supported by the Natural Sciences and Engineering Research Council of Canada.

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Publisher Item Identifier S 0018-9480(02)03022-3.

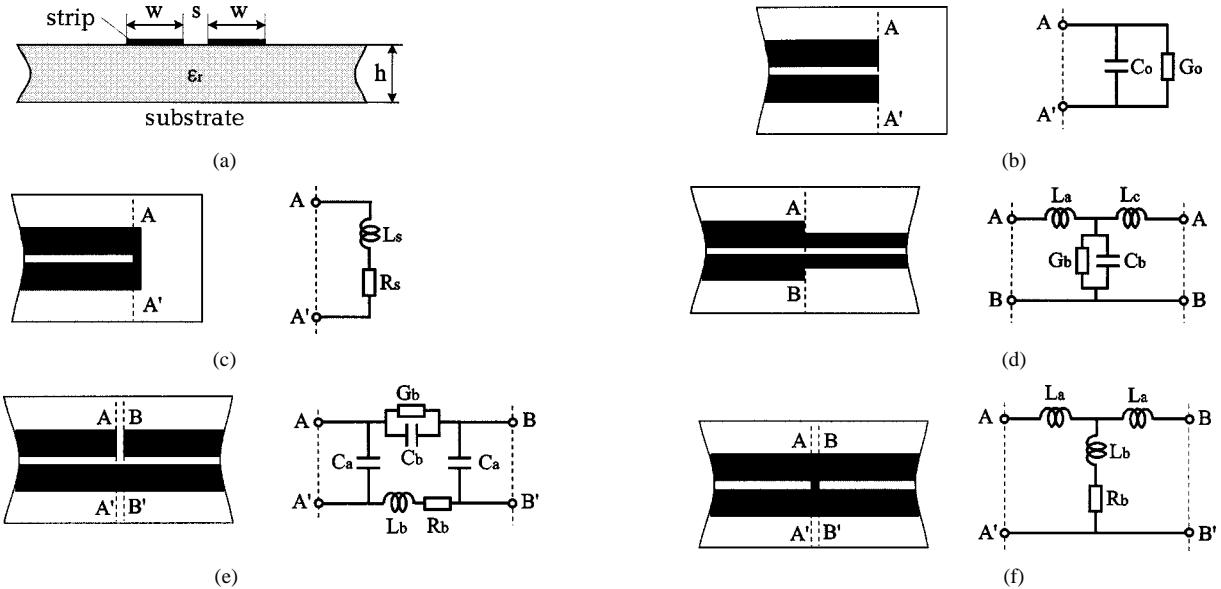


Fig. 1. Geometrical description, layout view, and equivalent-circuit model of CPS circuits and discontinuities. (a) Cross section of a uniform line. (b) Open circuit. (c) Short circuit. (d) Step discontinuity. (e) Capacitively coupled circuit (gap). (f) Inductively coupled circuit.

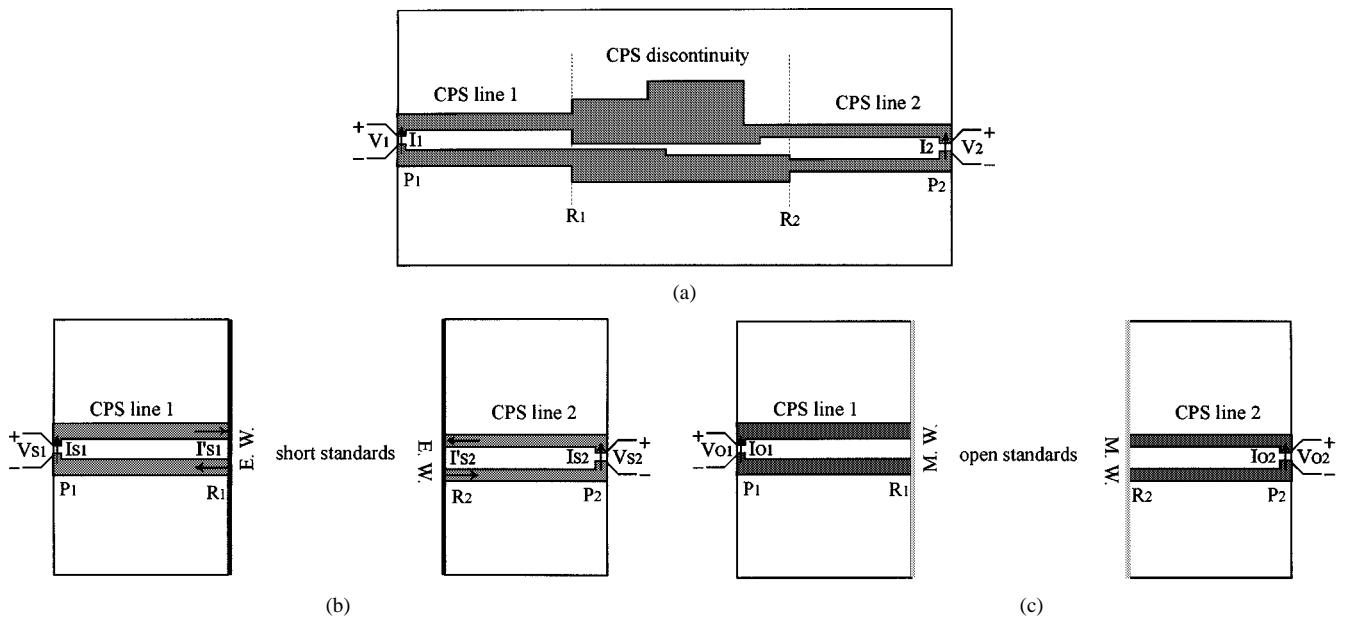


Fig. 2. Physical layout for the MoM characterization of a general CPS discontinuity and algorithmic arrangement of the two standard elements for the SOC parametric extraction of models from the MoM calculations. (a) Discontinuity layout. (b) Short standard element. (c) Open standard element.

technique, a very accurate equivalent-circuit model of a CPS circuit or discontinuity can be extracted from its physical layout, which accounts for all the physical effects such as frequency dispersion, high-order modes, and radiation loss. Interesting and unique features of CPS structures are described for the first time with their equivalent circuit. Comparison is also made with available data of some simple CPS circuits.

## II. FIELD MODELING AND PARAMETRIC EXTRACTION OF CPS CIRCUITS

Fig. 1(a) describes a uniform CPS with its cross-sectional view, and the two strip conductors with identical width are printed on the top interface of a dielectric substrate. Fig. 1(b)–(f)

show the top view of selected CPS discontinuities and their related lumped-element circuit models, namely, open, short, step, capacitively, and inductively coupled circuits. Note that it is not necessary to know *a priori* the equivalent lumped-element network since the SOC technique will extract the parameters in a self-consistent manner. In the following, a short narrative description will be given for the algorithm.

### A. Deterministic MoM Algorithm

Fig. 2(a) gives the geometry of a generalized CPS discontinuity for the MoM characterization, which is externally connected with two different CPS lines. One of the difficulties in the development of a deterministic MoM algorithm is how to effectively represent the impressed fields along the external feed lines

having different widths, which is usually called “the excitation mechanism” [18]. To explicitly derive the circuit parameters from our MoM calculation, an impressed electric field is introduced between the two strip conductors of two individual CPS lines at adequate planes of reference. As in Fig. 2(a), the port topology is formulated by longitudinally truncating and transversely extending the two strip conductors. Of course, the port locations are chosen electrically far away from the CPS discontinuity such that the parasitic interaction among the fields of the ports and the discontinuity is avoided.

As described in [11], an electric-field integral equation (EFIE) governing the electric current density ( $\vec{J}$ ) over the entire strip surface of the CPS lines and discontinuity can be derived in relation to the impressed electric field ( $\vec{E}^{\text{im}}$ ) at the port planes. Applying the Galerkin’s technique leads to a source-type matrix equation, allowing the calculation of the unknown current density distributed at each port location and CPS lines, as well as the discontinuity. To involve a numerical calibration procedure in our MoM algorithm, the entire CPS circuit is further classified into two parts, separated by two additional reference planes ( $R_1$  and  $R_2$ ): two uniform line blocks in which propagates only a dominant mode excited by the voltage launcher, and discontinuity region to which all the discontinuity effects are attributed. This geometrical arrangement is inspired by the basic concept of the well-known experimental calibration procedure, e.g., [14], namely, error boxes and device-under-test (DUT).

### B. SOC Calibration Scheme

The proposed SOC technique is then integrated with our MoM algorithm described above. It allows the parametric extraction of the core part of the CPS discontinuity (DUT) by evaluating and removing the two error boxes, which account for the port discontinuity and 3-D behavior of a uniform CPS line section for each external line. These error terms may bring about a “numerical noise” if a direct deembedding procedure is implemented that only relies on the 2-D behavior of a uniform CPS line section such as the characteristic impedance and electrical length. Even though very small, these terms are usually frequency or geometry dependent, as explained in [15]. Thus, they may be extremely harmful in the parameter extraction, considering that the CPS discontinuities of Fig. 1 are electrically small. This is because the lumped-element values of these discontinuities may become comparable with those of the error terms in this case.

The SOC technique requires two standard CPS short and open elements, as shown in Fig. 2(b) and (c), in connection with the two external lines of the CPS discontinuity. They are used to evaluate electrical behaviors of each CPS line section between the port location and the reference plane of the discontinuity, as depicted in Fig. 2(a). In the algorithm, the two standard elements can exactly be formulated by exciting a pair of odd and even impressed voltages at the two ports of a uniform CPS line with its length twice that of the corresponding CPS line section. In this way, the central location of the line can ideally be simulated by perfect electric (short-end) or magnetic (open-end) wall. Therefore, the uniform CPS line is divided into two identical parts, one of which represents the geometrical layout of the

CPS short or open element. Via a 3-D MoM characterization of the two CPS elements, the current profile over the entire strip surface can be calculated. As the currents flowing at the port planes of the two CPS standard elements are derived, the error terms (boxes) between the ports and reference planes of the CPS discontinuity can analytically be formulated through the solution of three groups of equation on the basis of network theory. Further, the characteristic  $ABCD$  matrix  $[X_i]$  can be derived in a closed form by using the port currents upon the two excitations, and the current flowing at the short end, namely,  $I_{si}$ ,  $I_{oi}$ , and  $I'_{si}$ , normalized with respect to the port voltages  $V_{si}$  and  $V_{oi}$  for the  $i$ th CPS external line

$$[X_i] = \begin{bmatrix} \bar{I}_{si} & -\frac{I}{\bar{I}'_{si}} \\ \frac{\bar{I}'_{si}}{\bar{I}'_{si} \bar{I}_{oi}} & \frac{\bar{I}'_{si}}{\bar{I}_{si} - \bar{I}_{oi}} \end{bmatrix} \quad (1)$$

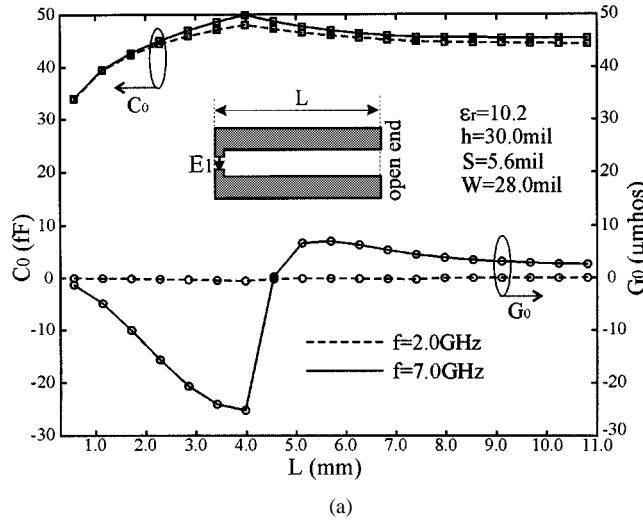
in which,  $\bar{I}_{si}$ ,  $\bar{I}_{oi}$ , and  $\bar{I}'_{si}$  represent the three normalized currents. Thus, the error terms related to the CPS lines can accurately be evaluated and removed from the calculated results.

## III. CPS CIRCUIT MODELS AND DISCUSSION

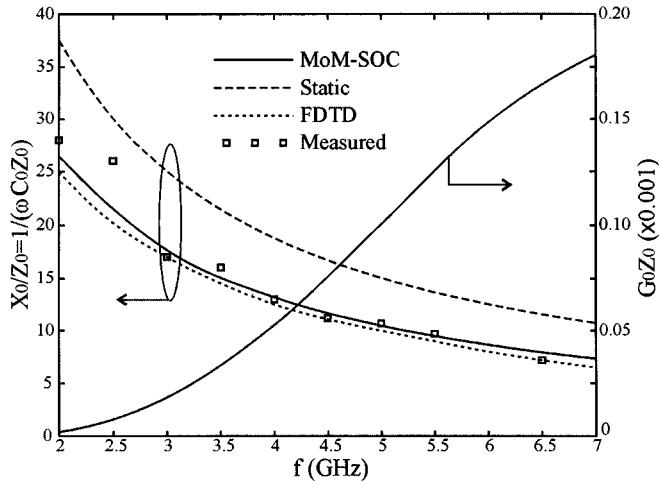
With the proposed SOC scheme in the MoM algorithm, a CPS circuit or discontinuity layout can be characterized as a general-purpose equivalent-circuit model with accurately extracted lumped elements. We begin with the investigation on the numerical stability and accuracy of our SOC scheme in the 3-D MoM by modeling of a simple CPS open circuit, prior to dealing with a variety of CPS circuits and discontinuities. Our special interests in this study are focused on the theoretical study of parasitic effects happening around the discontinuity of unbalanced CPS circuits, which are essentially generated by the finite strip width of the CPS lines.

### A. Open Circuit

Let us look at the circuit model and electrical characteristics of a CPS open circuit, as depicted in Fig. 1(b). Strictly speaking, the open end should equivalently be modeled as a shunt capacitance  $C_0$  and conductance  $G_0$ , standing for the fringing and radiating effects. These circuit parameters have been studied to some extent with a quasi-TEM model [19], thru-reflection line (TRL) measurements, and FDTD [6] as well as spectral-domain methods [9]. Unfortunately, there exist some intolerable differences among them and also strong frequency-dependent (or distributed) characteristics of all the deembedded shunt capacitance  $C_0$ . To our knowledge, it is basically attributed to the fact in the deembedding procedures that all the parasitic effects at the incident or input port in the numerical technique [6], [9] are completely ignored while the short-end element defined in the TRL measurements [6] is not a perfect reflection standard. Following the description in [11], they are very harmful in the parametric extraction of an electrically small CPS open circuit in such a way that the extracted  $C_0$  strongly depends on the operating frequency and external line length of choice. To validate the effectiveness of any deembedding procedure, one of



(a)



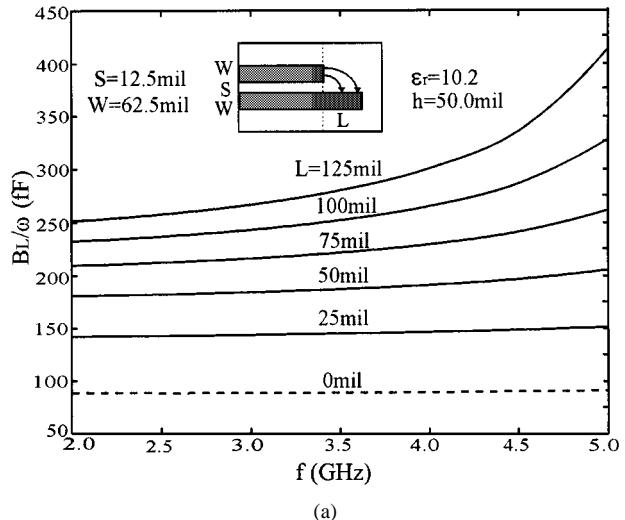
(b)

Fig. 3. Extracted complex admittance ( $Y_o = G_o + jB_o$ ) of a CPS open circuit. (a) Numerical convergence versus the feed-line length ( $L$ ). (b) Comparison with other theoretical and experimental results.

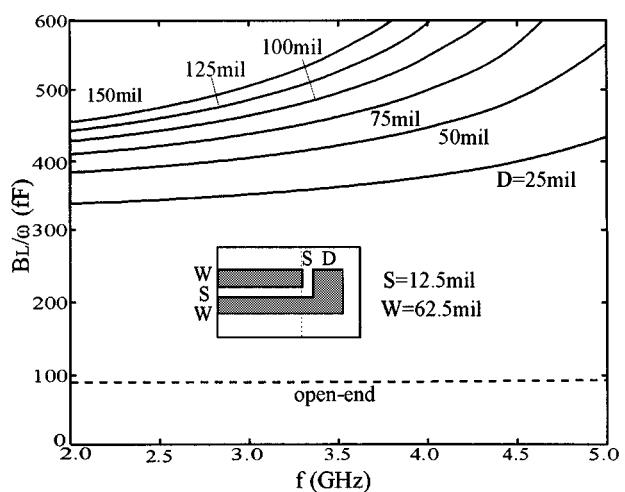
the standard procedures is to check whether the extracted circuit parameters are independent of the different external CPS lines ( $L$ ).

Fig. 3(a) depicts our extracted capacitance ( $C_0$ ) and conductance ( $G_0$ ) of a simple CPS open circuit against the CPS feed-line length ( $L$ ) at frequencies  $f = 2.0$  and  $7.0$  GHz. For the length ( $L$ ) shorter than  $7.0$  mm, both  $C_0$  and  $G_0$  appears significantly related to  $L$ , especially for  $G_0$  in the case of  $f = 7.0$  GHz. It may be attributed to a parasitic interaction of fields at the open-end and the port location over this short line range. As  $L$  is extended beyond  $7.0$  mm, they gradually converge to some stable values, representing their relative lumped-element parameters at the two specific frequencies. This result indicates that the problem of length-dependent instability can completely be solved with our SOC scheme if the port location is chosen far away from the core discontinuity.

Fig. 3(b) gives a general comparison of our extracted parameters against the other results over the frequency range of interest for the normalized reactance  $X_0/Z_0$ , as in [6]. Our  $X_0/Z_0$  is much closer to the FDTD results than the static and measured results over the frequency range from  $2.0$  to  $7.0$  GHz. Further,



(a)



(b)

Fig. 4. Parametric behavior of the extracted susceptance for two nonstandard CPS open circuits with extended line (called "type-A") and bent-line 90° section (called "type-B") as a function of frequency. (a) Type-A. (b) Type-B.

the extracted  $C_0$  is found well stationary at the value of  $44.8$  fF within a variation of  $1.5\%$ , thus exhibiting the lumped-element behavior of the CPS open circuit. The extracted  $G_0$  shows an exponential increase and then a linear increment with  $f$ , but its value is extremely smaller, as compared to its real counterpart  $B_0$  or  $1/X_0$ .

As compared with the infinitely extended ground plane of the other planar lines such as microstrip and CPW, the two strip widths of a CPS line are electrically narrow, susceptibly giving rise to a nonnegligible influence on electrical behaviors of some CPS circuits. In [6], a CPS circuit with an extended strip conductor was modeled as a lumped capacitance and it was found to be equal to three or four times of its conventional counterpart for a selected dimension. However, no detailed discussion was given as to its operating principle and physical insight. Judging from our extracted lumped elements of an open circuit (called "type-A") defined at its reference plane (open terminal), as shown in Fig. 4(a), we observe that the normalized susceptance  $BL/\omega$  exhibits an unexpectedly nonlinear behavior with frequency. This is much more pronounced with a large value of  $BL/\omega$  at higher frequency if one of the conductor strips is ex-

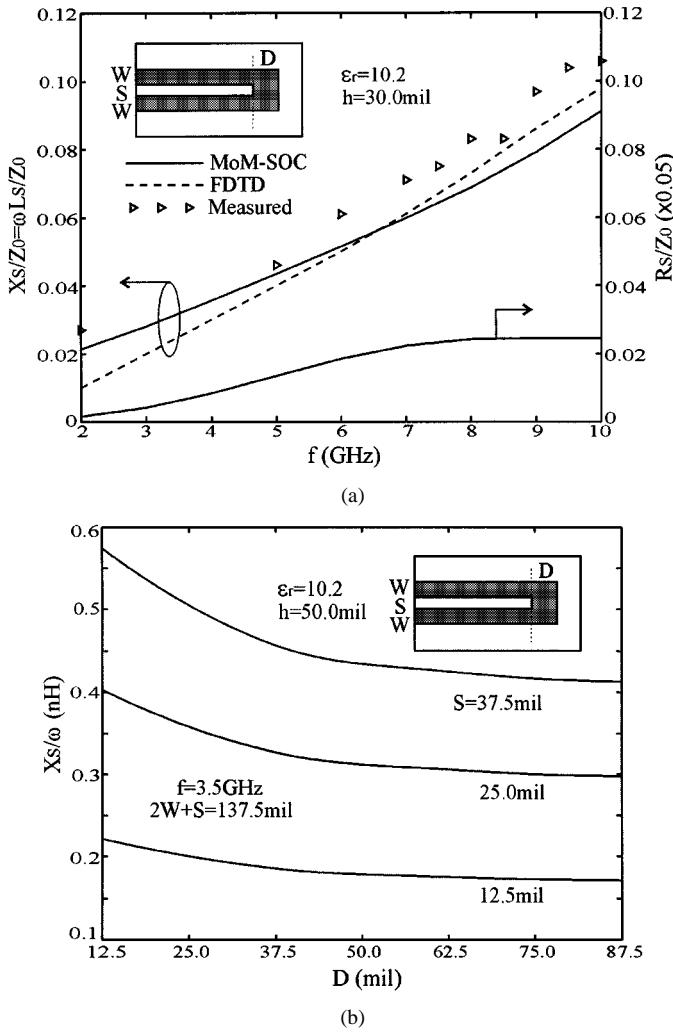


Fig. 5. Extracted complex short-end impedance ( $Z_s = R_s + jX_s$ ) of a CPS short circuit. (a) Comparison with the FDTD and experimental results. (b) Reactance ( $X_s/\omega$ ) versus the width ( $D$ ).

tended largely beyond the other with a length of  $L$ . Quite similarly, such a nonlinear property of dispersion can also be found in the extracted model of the other CPS open circuit (called “type-B”), as indicated in Fig. 4(b), in which one of the strip conductors is bent around the other. This interesting phenomenon can be explained based on a distributed coupling theory. In fact, a fringing-field coupling takes place between the open end of the short line and the extended line. It can be expected that this coupling is frequency dependent. In this case, the coupled section can equivalently be described by a circuit model in terms of a group of distributed shunt capacitances and series inductances rather than a purely lumped-capacitance model, as in Fig. 1(b). In general, these series inductances are caused from the extended conductor section having a finite strip width.

### B. Short Circuit

A CPS short circuit is realized by simply interconnecting the ends of two strip conductors, which can be described by a circuit model in terms of a series-connected inductance ( $L_S$ ) and resistance ( $R_S$ ), as shown in Fig. 1(c). Fig. 5(a) displays our extracted normalized reactance ( $X_S/Z_0$ ) and resistance ( $R_S/Z_0$ ) as a function of frequency against the FDTD and experimental

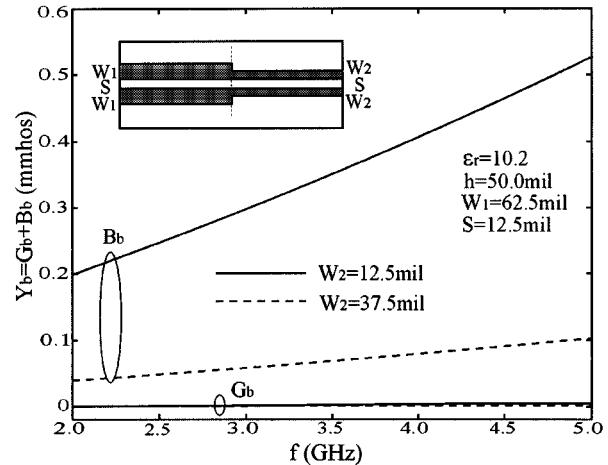


Fig. 6. Frequency-dependent shunt admittance ( $Y_b = G_b + jB_b$ ) of a CPS step discontinuity that is formed by two interconnecting strip widths ( $W_1$  and  $W_2$ ).

results [6]. Our  $X_S/Z_0$  increases in an approximately linear function with frequency and agrees reasonably with the other results over a wide frequency range (2.0–10.0 GHz). Otherwise,  $R_S/Z_0$  gradually rises up with frequency, representing the radiation loss from the transverse strip conductor with the width ( $D$ ) at the CPS short-end. Fig. 5(b) shows the extracted  $X_S/Z_0$  as a function of the strip width ( $D$ ) at  $f = 3.5$  GHz. It is found that  $X_S/Z_0$  starts to increase and becomes more pronounced for a wide slot width ( $S$ ) once the width ( $D$ ) is chosen smaller than 50.0 mil. To form a better CPS short circuit, a preferred way is to choose the width ( $D$ ) as wide as possible in reducing such an inductive effect of the short end. However, such a CPS short circuit definitely represents a nonnegligible inductance at the short end due to the finitely extended slot width ( $S$ ) so that this circuit element cannot be perceived as a perfect reflection standard in the deembedding or calibration procedure [6].

### C. Step Discontinuity

Similar to other planar structures, a CPS step discontinuity stands for an abruptly longitudinal change in linewidth, as described in Fig. 1(d), and its circuit model is generally represented with four lumped elements. In our calculations, only the complex shunt admittance ( $Y_b = G_b + jB_b$ ) is found to play a principal role in the entire equivalent network of such a CPS discontinuity. Fig. 6 gives our extracted shunt susceptance  $B_b$  and conductance  $G_b$  as a function of frequency for a CPS step discontinuity with a fixed slot width ( $S$ ). In this case,  $B_b$  goes up as a good linear function of frequency, thus exhibiting an excellent lumped-capacitance characteristic. Also, it shifts up rapidly with the ratio of two different strip widths of this CPS step, e.g.,  $W_1/W_2$ . On the other hand,  $G_b$  is found to increase visibly with frequency and width ratio ( $W_1/W_2$ ), but it still seems negligibly small as compared to  $B_b$  over the frequency range.

### D. Inductively Coupled Circuit

Fig. 1(f) describes a physical layout and its equivalent T-type circuit model of an inductively coupled CPS circuit with a transverse strip conductor bridging the two strip conductors. The equivalent network consists of four lumped-parameters, namely,

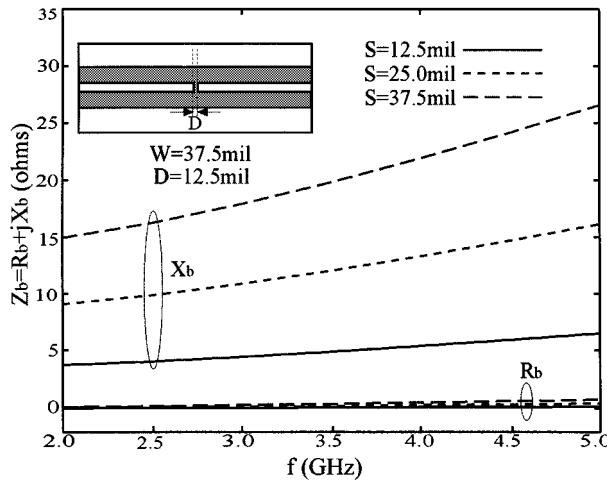


Fig. 7. Frequency-dependent shunt impedance ( $Z_b = R_b + jX_b$ ) of a CPS inductively coupled circuit with different CPW slot width ( $S$ ).

self-resistance/mutual resistance and reactances  $R_a$ ,  $R_b$ ,  $X_a$ , and  $X_b$ . These parameters can effectively be extracted from the full-wave MoM via our SOC scheme. Fig. 7 depicts the extracted mutual  $R_b$  and  $X_b$  as a function of frequency versus different slot width ( $S$ ) for a fixed narrow width ( $D = 12.5\text{ mil}$ ). It can be observed that  $X_b$  gradually increases with frequency ( $f$ ) and rapidly shifts up as the slot width ( $S$ ) is widened from 12.5 to 37.5 mil, exhibiting a great enhancement of its quasi-lumped series inductance ( $L_b$ ). In parallel,  $R_b$  slightly rises up as  $f$  increases and  $S$  is widened, but it is smaller than  $X_b$ . On the other hand,  $R_a$  and  $X_a$  is also much small in the case of the narrow  $D$  and wide  $S$  so that this CPS circuit can well be modeled by a single shunt reactance  $X_b$ .

#### E. Capacitively Coupled Circuit

As our last example, a capacitively coupled circuit presents some interesting and distinct features of the CPS circuits. This circuit consists of an air gap along one of the strip conductors while the other line is kept intact and this structure was also called a “CPS series gap,” as in [6]. Fig. 8(a) shows its geometrical view with a gap interval  $D$ . In [6], this structure was modeled as a lumped-capacitance  $\pi$ -network, and its FDTD-calculated parameters were plotted against measured ones at low frequency. Strictly speaking, its generalized circuit model should be considered as in Fig. 8(b) by adding a series inductance because the strip conductors have inherently electrically finite widths. This nonignorable parameter represents the parasitic effect of a noninfinitely extended strip width around the air gap section. The resulting circuit model of Fig. 8(b) can be converted into a  $\pi$ -network, as shown in Fig. 8(c), in which the series capacitance and inductance are put together to make up an equivalent series resonant circuit. In our MoM algorithm, any potential radiation loss [11]–[13] can be taken into account and can equivalently be incorporated into the model by a shunt conductance and a series resistance, as indicated in Fig. 8(b) and (c).

Fig. 9 shows our extracted circuit parameters over a wide frequency range (2.0–10.0 GHz) with the same dimensions as those in [6]. From Fig. 9(a), the self-susceptance ( $B_p$ ) is found

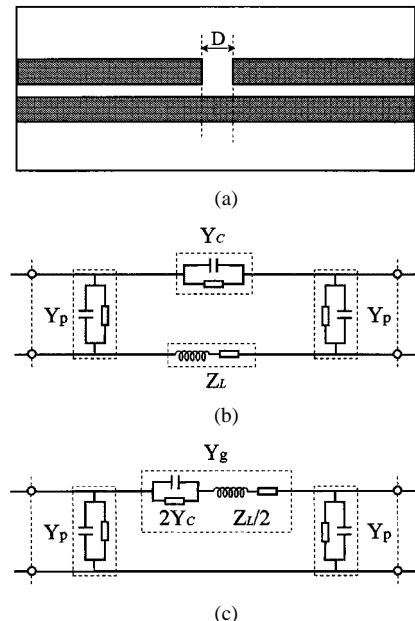


Fig. 8. Geometrical description and its possible equivalent circuit models of a CPS capacitively coupled (gap) circuit. (a) Top-view layout. (b) Original circuit model. (c) Convenient circuit model.

to linearly increase with frequency, while the self-conductance ( $G_p$ ) appears unchanged and much smaller than  $B_p$ . Fig. 9(b) gives the frequency response of the mutual-conductance ( $G_g$ ) and the mutual-susceptance ( $B_g$ ), depicting a vivid behavior of some unexpected radiating resonance, which is caused by the simultaneous existence of a series gap capacitance and a series inductance. Usually, the capacitance is intuitively obtainable by static or full-wave solutions, whereas the inductance due to the grounding line or the line of reference may easily be overlooked in the modeling. Generally speaking, a large gap interval triggers a large inductance under the fixed  $W$  and  $S$ , such that the resonant frequency is significantly reduced because the capacitance ( $C$ ) remains unchanged. Of course, the maximum radiation takes place around the frequency of resonance, translating the maximum value of  $G_g$ .

To validate our claimed results here, Fig. 10 gives a general comparison between our predicted results and those of [6], which depicts a frequency response of  $S$ -parameter characteristics for a structure with a wide strip width. Note that only low-frequency results are available in [6]. From our obtained curves, we can observe that a sharp variation of the  $S$ -parameters within the high portion of the frequency band is observed, which was not reported in [6]. This corresponds to a circuit resonance around  $f = 7.4\text{ GHz}$ , caused by a series gap capacitance and a series line inductance shown in Fig. 8(b) and (c). Interestingly, our recent work [20] demonstrates from both theory and experiments that such a resonant phenomenon exists in a finite-ground CPW gap circuit that also has an electrically finite-ground width. To our knowledge, there is a common physical mechanism for them that the finite ground at the same location as the gap interval ( $D$ ) generates an additional series inductance in the complete model, as illustrated in Fig. 8(b).

Now, let us take a close look at some visible and nonnegligible difference among three sets of the independent results in

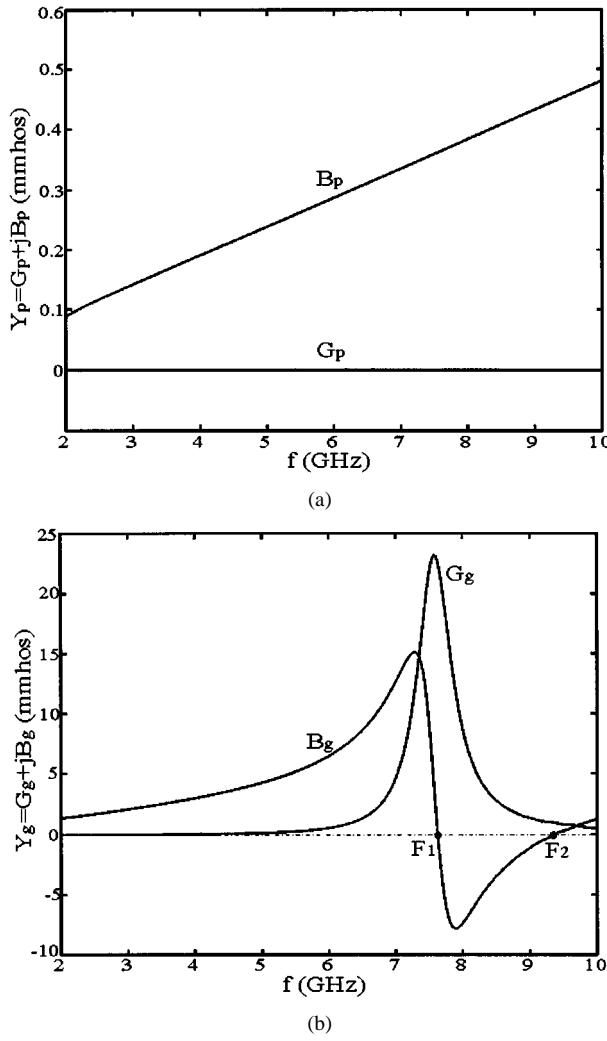


Fig. 9. Extracted complete circuit parameters of the CPS inductively coupled circuit. (a) Complex self-admittance ( $Y_p = G_p + jB_p$ ). (b) Complex mutual-admittance ( $Y_g = G_g + jB_g$ ).

both magnitudes ( $|S_{11}|$  &  $|S_{21}|$ ) and phases ( $\Phi_{11}$  &  $\Phi_{21}$ ) of  $S$ -parameters, as illustrated in Fig. 10. According to the principle of a reciprocal, lossless, and symmetrical two-port network [21], these magnitude and phase parameters should rigorously be satisfied with the following two basic equations, i.e.,  $|S_{11}|^2 + |S_{21}|^2 = 1$  and  $(\Phi_{11} - \Phi_{21}) = (2n - 1)90^\circ$ , where  $n$  is an integer number. Fig. 11 indicates the three sets of frequency-dependent loss factor ( $1 - |S_{11}|^2 - |S_{21}|^2$ ) and phase factor ( $\Phi_{11} - \Phi_{21}$ ), which are directly obtained from the FDTD and experimental results, as well as our SOC results as in Fig. 10. Over the frequency range ( $f < 4.8$  GHz) regarding the case of an extremely low loss [6], our results show that the loss factor is lower than 0.02, while the phase factor is almost kept around  $90^\circ$ , exhibiting an excellent consistency with the condition of a lossless two-port network. Unfortunately, the FDTD and experimental results illustrate a high loss factor with some irregularly frequency dependency even at low frequency. On the other hand, the phase factor for the FDTD results is almost kept around  $75^\circ$  far away from  $90^\circ$  in the lossless case, while randomly changes between  $60^\circ$  and  $100^\circ$  for the measured ones. To our understanding, it is mainly attributed to the unsolved issues such as nonideal calibration standards in the measurement and the com-

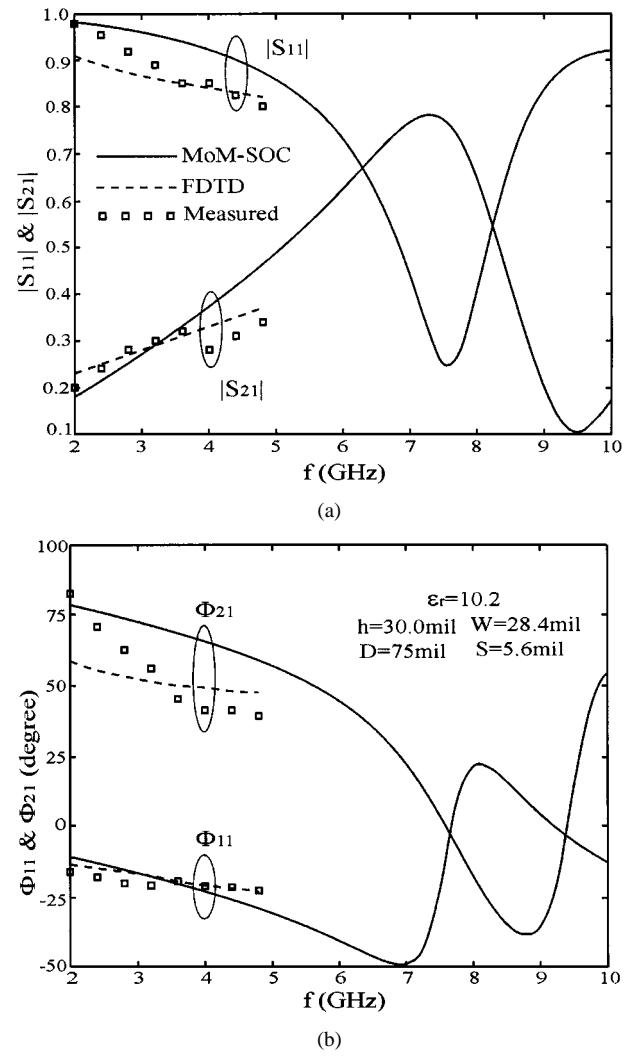


Fig. 10. Comparison of our calculated  $S$ -parameters against the FDTD and experimental results for a CPS capacitively coupled (gap) circuit, for which our results show an unexpected resonance around 7.6 GHz. (a) Magnitude of  $S$ -parameters:  $|S_{11}|$  &  $|S_{21}|$ . (b) Phase of  $S$ -parameters:  $\Phi_{11}$  &  $\Phi_{21}$ .

plete ignorance of the incident port discontinuity in the FDTD modeling for the parametric extraction of electrically small CPS circuits by using the existing deembedding procedures.

Furthermore, Fig. 11(a) illustrates a detailed radiation behavior of such a CPS gap discontinuity over a wide frequency range, in which the maximum radiation is observed to happen around the resonance with reference to Fig. 9(b). Fig. 11(b) indicates that the phase factor gradually goes down from the  $90^\circ$  line, drops off toward a negative minimum degree at a frequency location marked by “ $F_1$ ,” and then rises up to a positive degree as frequency is further increased beyond the location “ $F_2$ .” With reference to the results in Fig. 9(b), we can further observe that the susceptance ( $B_g$ ) is exactly equal to zero value at these two frequencies ( $F_1$  and  $F_2$ ) and becomes positive or negative values before or beyond them due to the capacitive or inductive  $B_g$ . Otherwise, we can observe that the phase factor is continuously varied as a function of frequency without any abrupt change, as in the lossless case regarding the dotted line in Fig. 11(b), and also its value is far away from the  $\pm 90^\circ$  lines, especially around the resonance. It can be well understood by the increasing radiation loss around the resonance, as shown in Fig. 9(b).

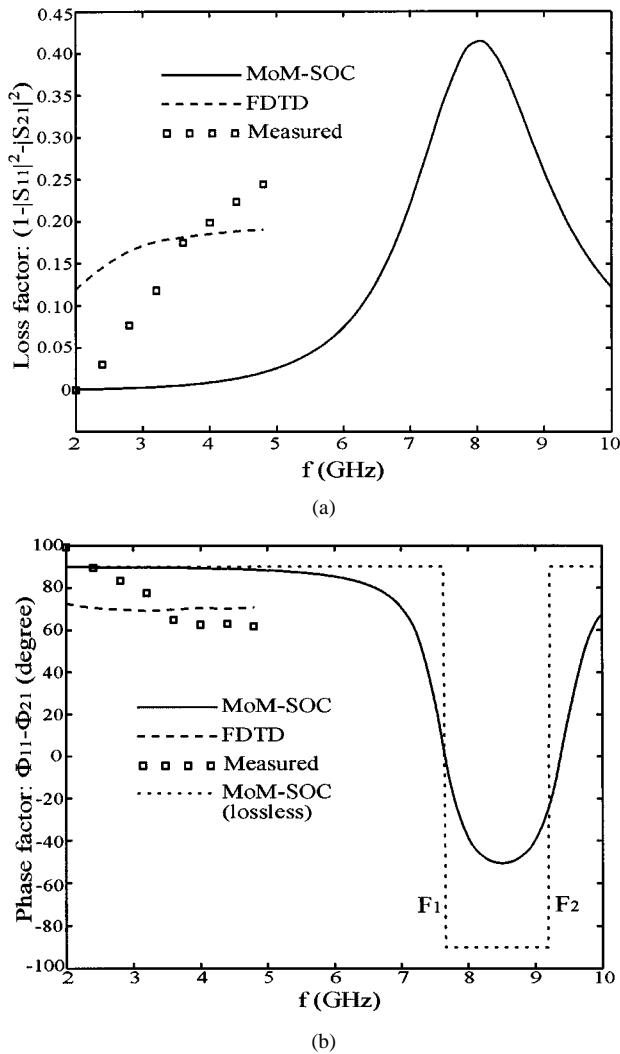


Fig. 11. Physical view of the three sets of  $S$ -parameters in Fig. 10 according to the principle of a reciprocal two-port network. (a) Loss factor:  $(1 - |S_{11}|^2 - |S_{21}|^2)$ . (b) Phase factor:  $(\Phi_{11} - \Phi_{21})$  (in degrees).

#### IV. CONCLUSION

In this paper, a variety of uniplanar CPS circuits and discontinuities have been studied in depth with the use of a deterministic 3-D MoM algorithm that incorporates a numerical deembedding technique called SOC. This unique scheme allows one to remove (calibrate) potential errors of the algorithm and extract the core circuit parameters in a very accurate manner. Thus, equivalent lumped-element circuit models can be obtained from physical layouts of the CPS discontinuities without resorting to a pre-knowledge or presetting of their circuit behavior. Interesting and distinct features of the CPS discontinuities and circuits are presented with detailed results of the extracted circuit models. It is discussed and explained for the first time that the finite-ground line of unbalanced CPS structures may generate nonnegligible and significant parasitic effects including an unexpected inductance effect or potential radiation loss or parasitic resonance phenomenon. Our CAD-oriented models provide a robust physical understanding and some insightful design rules of this new transmission-line-based circuits and structures, which is emerging as one of the new-generation building blocks of high-density radio-frequency design.

#### ACKNOWLEDGMENT

The authors gratefully acknowledge the anonymous reviewers for their valuable comments, which enabled the authors to produce a more accurate and readable revised manuscript.

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